Variable Interstellar Absorption toward the Halo Star HD 219188 — Implications for Small-Scale Interstellar Structure ¹

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ABSTRACT

Within the last 10 years, strong, narrow Na I absorption has appeared at $v_{\odot} \sim -38$ km s⁻¹ toward the halo star HD 219188; that absorption has continued to strengthen, by a factor 2–3, over the past three years. The line of sight appears to be moving into/through a relatively cold, quiescent intermediate velocity (IV) cloud, due to the 13 mas/yr proper motion of HD 219188; the variations in Na I probe length scales of 2–38 AU/yr. UV spectra obtained with the *HST* GHRS in 1994–1995 suggest $N(\rm H_{tot}) \sim 4.8 \times 10^{17}$ cm⁻², "halo cloud" depletions, $n_{\rm H} \sim 25$ cm⁻³, and $n_e \sim 0.85$ –6.2 cm⁻³ (if $T \sim 100$ K) for the portion of the IV cloud sampled at that time. The relatively high fractional ionization, $n_e/n_{\rm H} \gtrsim 0.034$, implies that hydrogen must be partially ionized. The $N(\rm Na~I)/N(\rm H_{tot})$ ratio is very high; in this case, the variations in Na I do not imply large local pressures or densities.

Subject headings: ISM: clouds — ISM: structure — Line: profiles — stars: individual (HD 219188)

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1. Introduction

Interstellar absorption lines were first recognized as interstellar by virtue of their invariance amidst the shifting stellar lines in the δ Ori A spectroscopic binary system (Hartmann 1904). Recently, however, it has become evident that some clearly interstellar lines do vary in strength and/or velocity. Variations in Na I and/or Ca II have been observed toward several stars in the direction of the Vela supernova remnant (Hobbs et al. 1991; Danks & Sembach 1995; Cha & Sembach 2000) — presumably due to cloud motions in that energetic environment. High N(Ca II)/N(Na I) ratios suggest significant shock processing of the interstellar dust grains in some of those clouds. In several other cases, however, the narrow Na I line widths ($b \sim 0.3$ –0.5 km s⁻¹) and/or low N(Ca II)/N(Na I) ratios characterizing the variable components suggest that the gas is relatively cool and quiescent (Blades et al. 1997; Price et al. 2000; Crawford et al. 2000; Lauroesch, Meyer, & Blades 2000). Detailed information regarding the physical conditions in these varying components is generally not available, however.

These indications of temporal variability may be related to recent results on small scale spatial structure in the ISM. Observations of H I 21cm absorption toward high-velocity pulsars have revealed structure on scales of 5–100 AU (Frail et al. 1994). Observations of absorption due to the trace ions Na I and/or K I toward globular clusters have found significant differences on scales of 10³-10⁶ AU (Langer, Prosser, & Sneden 1990; Kemp, Bates, & Lyons 1993; Meyer & Lauroesch 1999). Differences in Na I and/or K I absorption toward the individual members of binary or multiple stellar systems are commonplace, suggesting ubiquitous structure on intermediate scales of 10^2 – 10^4 AU (Watson & Meyer 1996; Lauroesch & Meyer 1999). If these differences are due to variations in overall hydrogen column density [assuming "typical" N(Na I)/N(H) ratios], the high implied densities (generally several thousand cm⁻³) are difficult to reconcile with clouds in thermal pressure equilibrium at typical interstellar pressures of order 3000 cm⁻³K. The differences may reflect a population of cold, dense filaments or sheets, embedded in warmer, less dense neutral gas and containing 10–30% of the total column density of cold, neutral gas (Heiles 1997). Alternatively, the differences in the trace ions might be due to variations in density and/or ionization on those small scales (e.g., Lauroesch et al. 1998). Differences in ionization would not explain the variations seen in HI, however.

HD 219188 is an early B supergiant located about 2800 pc from the sun at $(l,b) \sim (83^{\circ}, -50^{\circ})$. Because of its location and brightness $(V \sim 6.9)$, it has been observed in a number of studies investigating the interstellar clouds in the lower Galactic halo (Münch & Zirin 1961; Albert 1983; Danly 1989; Sembach, Danks, & Savage 1993). Strong, multi-component absorption from Na I and Ca II is apparent for -17 km s⁻¹ $\lesssim v_{\odot} \lesssim +5$ km s⁻¹; weaker Ca II absorption extends to about -35 km s⁻¹. High resolution Na I spectra obtained in 1997, however, revealed a new, relatively strong, narrow component at $v \sim -38$ km s⁻¹. This component is also present in Ca II and in several species in UV (GHRS echelle) spectra obtained in 1994–1995. Subsequent monitoring of the line of sight shows that the column densities of both Na I and Ca II have increased by additional factors of 2–3 since 1997. In the following sections, we discuss the recent optical and UV spectra

of HD 219188 and the inferred properties of the newly-revealed intermediate-velocity cloud along that line of sight.

2. Optical Data

The initial Na I and Ca II spectra of HD 219188 were obtained with the coudé feed and coudé spectrograph at the Kitt Peak National Observatory, as part of a program to acquire high resolution (FWHM $\lesssim 1.5~\rm km~s^{-1}$) optical spectra of halo stars previously observed at slightly lower resolution in the UV with the *HST* GHRS (see Fitzpatrick & Spitzer 1997). After the 1997 Na I spectrum of HD 219188 revealed the new intermediate-velocity (IV) component at $-38~\rm km~s^{-1}$, additional spectra were obtained during various runs at Kitt Peak, the Anglo-Australian Observatory, and the European Southern Observatory in order to monitor possible further variations (Table 1).

The spectra were processed and extracted using standard procedures within IRAF, then normalized via polynomial fits to continuum regions adjacent to the absorption lines. The resulting line profiles were fitted with multiple Voigt profile components, yielding column densities, line widths, and heliocentric velocities for the individual components (e.g., Welty, Morton, & Hobbs 1996). Some of the Na I and Ca II profiles are shown in Figure 1, with the earlier profiles described by Albert (1983) and Sembach et al. (1993) at the top for comparison.

The changes in the strength of the -38 km s^{-1} component — over the past 20 years and over intervals as short as 6 months — are quite evident and striking (Fig. 1; Table 1). While weak Na I absorption at $v \sim -38 \text{ km s}^{-1}$ appears to have been present in 1991 (Sembach et al. 1993), none was detected in 1980 (Albert 1983). By 1997, the Na I column density was roughly a factor of 10 higher than the limit in Albert (1983), and it has increased by a further factor of 2 since then. As corresponding IV Ca II absorption may be present in the profiles from both earlier studies (Fig. 1), we have fitted those profiles using the component velocity structure derived from our higher resolution spectra to estimate N(Ca II) in the IV cloud at those earlier epochs. While our Ca II column densities in 1995 and 1997 may be consistent with the (uncertain) value reported by Albert (1983), N(Ca II) has increased by a factor of 2–3 since then. Over the past 3–5 years, the N(Na I)/N(Ca II) ratio thus appears to have remained roughly constant (though significantly higher than previous values or limits) (Table 1).

In contrast to the changes in column density, neither the velocity of the IV component nor the Na I line width has changed significantly since 1995. The width of the Na I line, $b \sim 0.55$ –0.60 km s⁻¹ implies a maximum temperature for the cloud of 490 K and a maximum 3-dimensional internal turbulent velocity $(3^{1/2})v_t = 0.73$ km s⁻¹ (comparable to the sound speed for $T \sim 80$ –100 K).

3. UV Data

HD 219188 was observed with the *HST* GHRS on 1994 June 6 and 1995 May 16–17, at resolutions of about 3.5 km s⁻¹, under GTO program 5093. Twelve wavelength settings were obtained, using the FP-SPLIT procedure to enable the identification and elimination of detector fixed-pattern features. The individual sub-exposures for each wavelength setting were combined using the POFFSETS and SPECALIGN routines within the STSDAS package. The summed spectra were normalized via polynomial fits to continuum regions adjacent to the absorption lines; the empirical S/Ns range from 30–50. The interstellar absorption line profiles were then fitted using the component structures obtained from the higher resolution spectra of Na I and Ca II (Fitzpatrick & Spitzer 1997; Welty et al. 1999).

The $-38~\rm km~s^{-1}$ component is detected in the ground state lines of C I, C II, O I, Si II, S II, and Fe II and in the excited fine-structure lines of C I*, C I**, and C II* (Table 2, Fig. 2). If sulfur is undepleted, then $N(\rm S~II)$ suggests a total IV hydrogen column density $N(\rm H_{tot}) = N(\rm H~I) + N(\rm H~II) \sim 4.8 \times 10^{17}~\rm cm^{-2}$. The IV column density estimated for O I (which should track H I) is consistent with that $N(\rm H_{tot})$ and the typical interstellar oxygen abundance $[\rm O/H] \sim -0.4~\rm dex$ (Meyer, Jura, & Cardelli 1998) — which suggests that the IV gas is largely neutral. The C II $\lambda 1334$ line is saturated, but we may use $N(\rm S~II)$ to estimate the total IV carbon abundance $[N(\rm C~I) + N(\rm C~II)] \sim 6.8 \pm 2.1 \times 10^{13}~\rm cm^{-2}$, as $[\rm C/S] \sim -0.4 \pm 0.1$ dex in diffuse Galactic clouds. The column densities for Fe II and Si II yield depletions similar to those found for clouds in the Galactic halo (Welty et al. 1999, and references therein).

Analysis of the C I fine-structure excitation equilibrium enables estimates for the thermal pressure and local density in the IV gas (Jenkins & Shaya 1979). The observed $N(\text{C I}^*)/N(\text{C I}) = 0.17\pm0.03$ and $N(\text{C I}^{**})/N(\text{C I}) = 0.06\pm0.02$ yield $n_{\rm H}T \sim 2500$ cm⁻³K and $n_{\rm H} \sim 25$ cm⁻³ for T = 100 K. The thickness of the region probed, $N(\text{H})/n_{\rm H}$, would thus be ~ 0.006 pc, or about 1300 AU. Because $n_{\rm H}T$ and $n_{\rm H}$ depend on the assumed T, we have listed values for several temperatures consistent with the limit from b(Na I) in Table 3.

Estimates for the electron density in the IV gas may be derived both from considerations of photoionization equilibrium and from C II excitation equilibrium. If photoionization equilibrium is assumed for carbon [with $\Gamma/\alpha=24$ for the WJ1 interstellar radiation field and T=100 K (Péquignot & Aldrovandi 1986) and with N(C II) as estimated above], then $N(\text{C I})/N(\text{C II})=n_e$ (Γ/α)⁻¹ yields $n_e(\text{C I})\sim 6.2\pm 2.4$ cm⁻³. Analysis of the C II fine-structure excitation (Fitzpatrick & Spitzer 1997; Welty et al. 1999) for T=100 K, however, yields $n_e(\text{C II})\sim 0.85^{+0.52}_{-0.46}$ cm⁻³—about a factor of 7 smaller. While the uncertainties in both n_e estimates are dominated by the almost 50% uncertainty in N(C II), the ratio of the two estimates is not very sensitive to N(C II).

Similar discrepancies between $n_e(C I)$ and $n_e(C II)$ have been noted for other lines of sight (Fitzpatrick & Spitzer 1997; Welty et al. 1999). While C I might be concentrated in a smaller, denser region than C II, $n_e(C I)$ for that smaller region would be even larger than the average value derived above, and would imply a fractional ionization n_e/n_H greater than 25%. The $n_e(C I)$

would be reduced, however, if the radiation field were weaker or if T were lower than assumed above. Alternatively, the large n_e (relative to the assumed WJ1 radiation field) suggests that enough negatively charged large molecules (e.g., PAH's) could be present for charge exchange reactions (X⁺ + LM⁻ - > X⁰ + LM⁰) to significantly enhance the column densities of the trace neutral species — in which case $n_e(C\ I)$ would be too high (Lepp & Dalgarno 1988; Lepp et al. 1988; Welty & Hobbs 2001). Even if that is the case, however, $n_e(C\ II)$ still yields $n_e/n_H \sim 0.034$ — more than a factor 240 greater than the value 1.4×10^{-4} that would be due to photoionization of carbon (usually assumed to be the primary source of electrons). Hydrogen therefore appears to be partially ($\gtrsim 3.4\%$) ionized in the -38 km s⁻¹ cloud. It is not clear how such a high fractional ionization could be maintained.

4. Discussion

The Leiden-Dwingeloo H I maps (Hartmann & Burton 1997) suggest that the line of sight to HD 219188 lies near the inner edge of a shell of IV gas, seen over the range -30 to -60 km s⁻¹ (v_{\odot}). Despite its possible association with an expanding and/or compressed shell, the cloud's fairly low T, fairly high N(Na I)/N(Ca II) ratio, moderate pressure, and stable velocity suggest that it is relatively quiescent. The variations in absorption may instead be due to the relatively high proper motion of HD 219188 — about 13.4 mas/year (ESA 1997) — which corresponds to a transverse velocity of about 180 km s⁻¹, or about 38 AU per year, at 2800 pc. The steadily increasing column densities of Na I and Ca II and the high fractional ionization suggest that the line of sight is moving into/through the partially ionized edge of the IV cloud due to the motion of the background star.

Unfortunately, the distance to this IV cloud is not well constrained. A lower limit of 150 pc is provided by the lack of IV absorption toward β Psc (<2°away, but toward detectable IV H I emission); the upper limit (2800 pc) is set by HD 219188 itself. The length scale in the IV cloud probed by the stellar proper motion is thus between 2 and 38 AU/year — considerably smaller than the 10^2 – 10^4 AU scales sampled by the binary star studies. The cloud distance (and thus the length scale) could be better constrained by obtaining Na I spectra of other stars in this region (at various distances).

While we do not have a spectrum of Na I from 1994–1995, when the UV spectra were obtained, estimates based on the C I and Ca II observed then suggest that N(Na I) could have been $\sim 2 \times 10^{11} \text{ cm}^{-2}$. The ratio $N(\text{Na I})/N(\text{H}) \sim 4.2 \times 10^{-7}$ would thus have been much higher than is typically observed for diffuse clouds with $N(\text{H}) \gtrsim 10^{19} \text{ cm}^{-2}$ (Welty & Hobbs 2001) — presumably due to the high fractional ionization. Furthermore, if Na is depleted by a factor 4 in the IV gas (Welty & Hobbs 2001), then Na I would have been the dominant form of Na.

If such high N(Na I)/N(H) ratios and fractional ionizations are representative of other sightlines where differences in Na I have been observed over small transverse scales, then the corresponding differences in N(H) — and the resulting inferred pressures and densities — would not be as large as previously estimated. Lauroesch et al. (2000) have estimated $n_e/n_{\rm H} \sim 0.008$ for the variable component at $v \sim 9$ km s⁻¹ toward HD 32040, and $N({\rm Na~I})/N({\rm H})$ is likely high [for $N({\rm H}) \sim 6 \times 10^{18}~{\rm cm}^{-2}$] in the -8.6 km s⁻¹ component toward μ^1 Cru (Meyer & Blades 1996). The high $N({\rm Na~I})/N({\rm H})$ ratios found for a sample of high- and intermediate-velocity clouds with $N({\rm H}) \lesssim 10^{19}~{\rm cm}^{-2}$ (Wakker & Mathis 2000) may also generally be due to partial ionization of the hydrogen. These lower $N({\rm H})$ clouds may thus represent a different regime from the more familiar diffuse neutral clouds with $N({\rm H}) \gtrsim 10^{19}~{\rm cm}^{-2}$.

Lauroesch et al. (1998) have conjectured that some of the small-scale structure observed in Na I may be due to fluctuations in ionization [rather than in overall $N({\rm H})$] for cases where corresponding differences are not seen for Zn II [which should track $N({\rm H})$]. The low overall column densities, moderate local densities, and high fractional ionization found for the IV cloud toward HD 219188 suggest that some of the small-scale structure seen in Na I and K I may be due to partially ionized, relatively low-density cloud edges and/or fragments — rather than dense, neutral clumps. Further monitoring of the HD 219188 line of sight over the next several years — particularly in the UV with HST/STIS — should provide interesting information on the variations of density, temperature, and depletions with depth in the IV cloud.

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Table 1. Optical Data: $v = -38 \ \mathrm{km \ s^{-1} \ Component^a}$

Date	Observatory	Resolution	Na I		Ca II			$N({ m Na~I})/N({ m Ca~II})$	
	(Instrument)		$N_{10}{}^{\rm b}$	b	v	N_{10}^{b}	b	v	
		$({\rm km~s^{-1}})$	$({\rm cm}^{-2})$	$({\rm km~s^{-1}})$	$({\rm km}\ {\rm s}^{-1})$	$({\rm cm}^{-2})$	$({\rm km~s^{-1}})$	$({\rm km}\ {\rm s}^{-1})$	
1980.75^{c}	McDonald (2.7m)	$5.9/5.4^{c}$	< 3.0			4 ± 1	(1.5)	(-38.6)	< 0.8
$1991.56^{ m d}$	ESO (CAT/CES)	4.4	4 ± 3	1.5 ± 0.9	-37.9 ± 1.1	2 ± 1	(1.5)	(-38.6)	$2.0 {\pm} 1.8$
1995.82	KPNO (coudé feed)	1.40				3.1 ± 0.3	(1.5)	-38.6 ± 0.2	
1997.77	KPNO (coudé feed)	1.35	27 ± 1	$0.59 {\pm} 0.06$	-38.3 ± 0.1	4.9 ± 1.0	(1.5)	-38.9 ± 0.3	5.5 ± 1.1
1998.68	AAO (AAT/UCLES)	5.0	34 ± 2	(0.6)	-38.0 ± 0.1		· ´		
1999.42	KPNO (coudé feed)	1.50	39 ± 4	(0.6)	-38.1 ± 0.1				
1999.98	ESO (3.6m/CES)	1.20	52 ± 2	$0.55 {\pm} 0.05$	-38.2 ± 0.1				
2000.46	KPNO (coudé feed)	1.50	60 ± 5	$0.56{\pm}0.09$	-38.3 ± 0.1	$10.7 {\pm} 1.2$	(0.6)	-38.3 ± 0.1	$5.6 {\pm} 0.8$

^aUncertainties are 1σ ; values in parentheses were fixed in the fits.

 $^{^{\}rm b}{\rm Column}$ density in units of $10^{10}~{\rm cm}^{-2}.$

 $^{^{\}rm c}{\rm Albert}$ 1983; FWHM for Na I/Ca II; $N({\rm Ca~II})$ from new fit to generated spectrum.

 $^{^{}m d}$ Sembach et al. 1993; shifted to align low-velocity absorption; $N({
m Ca~II})$ from new fit to generated spectrum.

Table 2. UV Data: $v = -38 \text{ km s}^{-1} \text{ Component}^{\text{a}}$

Ion	Line (Å)	$N(X)$ (cm^{-2})	[X/S] ^b
СІ	1560	$10.7 \pm 1.0 \mathrm{e} 12$	
$C I^*$	1560	$2.3 {\pm} 0.3 {\mathrm e} 12$	
C I**	1561	$0.8 {\pm} 0.3 {\mathrm e} 12$	
C II	1334	$(4.9\pm2.1e13)^{c}$	$(-0.4)^{c}$
C II*	1335	$5\pm1e12$	
ΟI	1302	$1.1 \pm 0.4 \mathrm{e} 14$	-0.5
Si II	1304	$9 \pm 1 e12$	-0.3
S II	1253	$9\pm2e12$	
Fe II	2374, 2600	$4.0 \pm 0.4 \mathrm{e} 12$	-0.6

^aUncertainties are 1σ ; limits are 2σ .

 $^{\rm b}[{\rm X/S}] = \log\{N({\rm X~II})/N({\rm S~II})\} - \log\{({\rm A_X/A_S})_{\odot}\},$ using the solar system meteoritic abundances of Anders & Grevesse 1989 and Grevesse & Noels 1993.

 $^{\rm c}Estimated$ assuming typical interstellar [C/S] = $-0.4{\pm}0.1$ dex (for C I + C II).

Table 3. Inferred Properties: $v = -38 \text{ km s}^{-1}$ Component

Quantity		Value	
$N({ m H})~({ m cm}^{-2})$		4.8×10^{17}	
T(K)		< 490	
	25 K	100 K	400 K
$\log(n_{\rm H}T)~({\rm cm}^{-3}{\rm K})$	3.4 ± 0.1	3.4 ± 0.1	3.8 ± 0.1
$n_{\rm H} \ ({\rm cm}^{-3})$	100	25	15
thickness (AU)	325	1300	2200
$n_e(\mathrm{C~I})^{\mathrm{a}}(\mathrm{cm}^{-3})$	2.6	6.2	14.7
$n_e(\mathrm{C~II})^\mathrm{b}(\mathrm{cm}^{-3})$		0.85	0.79
$n_e/n_{ m H}{}^{ m c}$		0.034	0.053

 $^{\mathrm{a}}\mathrm{Assuming}$ photoionization equilibrium for carbon; WJ1 radiation field.

 $^{\rm b}{\rm From~C}$ II fine-structure excitation equilibrium; no valid value for $T=25~{\rm K}.$

 $^{\rm c}{\rm Using}~n_e$ from analysis of C II fine-structure.

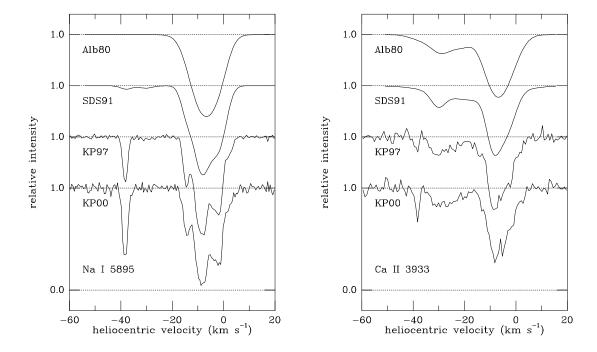


Fig. 1.— Selected Na I $\lambda 5895$ (left) and Ca II $\lambda 3933$ (right) spectra toward HD 219188. The sources and dates of the spectra are indicated; see Table 1. The column densities in the intermediate-velocity component at $v \sim -38$ km s⁻¹ have increased by factors of 2–3 over the past 3–5 years; no Na I was detected at that velocity in 1980. The lower resolution spectra from Albert (1983; Alb80) and Sembach et al. (1993; SDS91) were generated from the component structures listed therein.

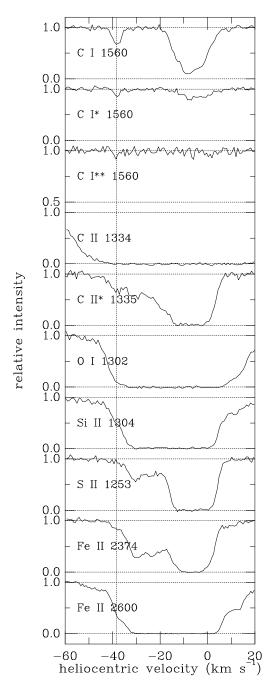


Fig. 2.— Profiles of selected UV lines toward HD 219188, observed with GHRS at resolutions of about 3.5 km s⁻¹ in 1994–1995. The vertical dotted line indicates the component at -38 km s^{-1} .